ISOTROPIC SYSTEMS
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The conductivity of multicomponent isotropic heterogeneous system is studied with the help of a computer using statistical analysis.

We will investigate an isotropic heterogeneous system consisting of different polyhedrons, whose components fill the entire space without voids and have different transfer coefficients $\Lambda_{i}$. If $i=2$, then a binary system is being examined, while for $\mathrm{i}>2$, a multicomponent heterogeneous system is being examined. In order to study the transfer coefficients of such materials, in recent years, computer-aided statistical techniques are widely used, which has permitted obtaining important characteristics, unified by the term theory of percolation or percolation processes. The general topological picture of the structure of a binary heterogeneous system as a function of the concentration of components has become clear from the work in [1-3]. Figure 1 represents a binary mixture of polyhedrons, where the dark regions have high conductivity, while the white regions are perfect insulators. Assume that initially the entire space is filled by insulators (Fig. 1a) and, then, a conductor is randomly disseminated into the system. A mathematical experiment on a computer, carried out using the Monte Carlo method, showed that for low concentration $\mathrm{m}_{\mathrm{M}}$ (M indicates a metal) of the conducting component, regions with high conductivity appear singly or as small clusters (Fig. 1b). As the system grows, large clusters can form as well, which together with the small clusters form so-called isolated clusters (IsC). When the concentration $m_{M}$ approaches a critical value $m_{C}\left(m_{M} \rightarrow m_{c}\right)$, the large clusters begin to coalesce with one another and giant clusters with odd shapes appear (Fig. 1c and f), separated from one another by a low nonconducting space $A-A$. For $m=m_{c}$, the isolated clusters coalesce and form infinite clusters (InC); the system becomes conducting. With further growth, $m_{M}>m_{C}$, the infinite cluster increases, absorbing small clusters, and conducting chains permeate the entire system, forming a system with interpenetrating components (Fig. 1d), and the latter, with $m_{M}=1$, becomes a homogeneous (metallic) system (Fig. 1e). The value $m=m_{c}$ is called the percolation threshold, at which, in a conductor-ideal insulator system, conductivity by "hops" increases from zero to some magnitude and then varies monotonically with increasing concentration $m_{M}$. This phenomenon is called hopping conductivity and the existence of a percolation thr eshold is a general phenomenon inherent both to disordered and lattice models. The effective conductivity $\Lambda$ of a two-component highly inhomogeneous system $\nu=\Lambda_{\mathrm{I}} / \Lambda_{\mathrm{M}}=$ 0 , consisting of a mixture of ideal insulators and conductors with conductivities $\Lambda_{I}$ and $\Lambda_{M}$, can be described with the help of the following computer-fitted equations [3, 4]:

$$
\begin{gathered}
m_{c}<m \leqslant 0,5, N / \Lambda_{\mathrm{M}}=A\left(m_{\mathrm{M}}-m_{\mathrm{c}}\right)^{K} \\
m<m_{\mathrm{c}}, N / \Lambda_{\mathrm{I}}=1 /\left(1-5 m_{\mathrm{M}}\right. \\
m_{\mathrm{c}}=0.15 \pm 0.03, K=1.8 \pm 0,2, A=1-1.6
\end{gathered}
$$

In what follows, the conductivity of inhomogereous systems was investigated with methods in which a computer statistical analysis was combined with different model representations. These methods are reviewed in [5, 6]. In particular, a model constructed based on a combination of the percolation method and reduction to an elementary cell is constructed in [6]. The equations obtained permit finding with satisfactory accuracy the conductivity of isotropic binary systems for any ratio of the conductivities of the components $0 \leq \Lambda_{1} / \Lambda_{2} \leq 1$ in the entire range of variation of their concentrations $0 \leq m_{1}=1-m_{2} \leq 1$.

The formation of an infinite cluster and the appearance of hopping conductivity have been studied quite well using statistical methods with the help of a computer for binary media. An important problem is the extension of this problem using the same methods to multicomponent (two and more components) heterogeneous

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Fig. 1. Model of a binary two-dimensional inhomogeneous body: a) $m_{M}=0 ; m_{I}=1 ;$ b) $m_{I}>m_{M}, m_{M}<m_{c} ; c$ and f) isolated clusters; d) $m_{M} \geq m_{c} ;$ e) $m_{M}=1, m_{I}=0$.


Fig. 2. Model of a heterogeneous system.
systems for arbitrary conductivity of the components.
We point out one more problem, which is most naturally solved by appealing to a mathematical experiment on a computer. In determining experimentally the thermophysical properties of inhomogeneous systems or in calculating their temperature fields, the inhomogeneous system is usually replaced by a quasihomogeneous body with some effective properties. In so doing, it is necessary to estimate the dimensions of the so-called representative element of the system, whose transport coefficient is approximately equal to the transport coefficient of the entire mass. In what follows, we present methods for solving such problems and results of solutions.

We examine first a binary heterogeneous system with a random distribution of components (Fig. 1) and we will choose from the system some cube, in which the concentration of components equals the corresponding concentration in the entire mass; we will assume that the lengths of the cube edges equal unity and that the lateral surfaces of the cube are adiabatic, i.e.,

$$
\begin{equation*}
\left.\frac{\partial t}{\partial x}\right|_{x_{x=0}}=\left.\frac{\partial t}{\partial x}\right|_{x=1}=0 ;\left.\frac{\partial t}{\partial y}\right|_{y=0}=\left.\frac{\partial t}{\partial y}\right|_{y=1}=0 \tag{1}
\end{equation*}
$$



Fig. 3. Dependence of the conductivity of a binary heterogeneous system on the number of layers in the model (a) [1) $\Lambda / \Lambda_{M}=1 \cdot 10^{-2}$, $\mathrm{m}_{2}=0.3$; 2) $1 \cdot 10^{-4}$ and 0.3 ; 3) $1 \cdot 10^{-2}$ and 0.7 ; 4) $1 \cdot 10^{-4}$ and 0.7 ; 5) $1 \cdot 10^{-2}$ and 0.9 ;6) $1 \cdot 10^{-4}$ and 0.9 ] and on concentration (b).


Fig. 4. Concentration dependence of the conductivity of a three-component system.
while the surfaces $\mathrm{z}=0$ and $\mathrm{z}=1$ are isothermal (Fig. 2)

$$
\begin{equation*}
t_{\left.\right|_{z=0}}=t^{\prime}, \quad t_{\left.\right|_{z=1}}=t^{\prime \prime} \tag{2}
\end{equation*}
$$

The basic method for studying the conductivity through such a system is to simulate the system with the help of a collection of connected resistances. We separate the cube being examined by planes parallel to the base into N layers and each layer in its turn into $\mathrm{N}^{2}$ parts with planes parallel to the lateral faces. Resistances of some type are positioned between the nodes of the lattice obtained and the number of each type of resistance is proportional to the concentration of the components (Fig. 2). With the help of a pseudorandom number generator, the magnitude of the resistances is assigned one of two possible values. We will denote the conductivities of the components by $\Lambda_{1}$ and $\Lambda_{2}$, and their concentrations by $m_{1}$ and $m_{2}\left(m_{1}+m_{2}=1\right)$, and we will find the overall resistance $R$ to a flow $Q$ between the isothermal surfaces:

$$
\begin{equation*}
R=\frac{t^{\prime \prime}-t^{\prime}}{Q} . \tag{3}
\end{equation*}
$$

For a single cube, the area of the isothermal surface and the height of the cube equal unity, so that the effective conductivity is $\Lambda=R^{-1}$. If we assume that $t^{\prime}=0$ and $t^{\prime \prime}=1$, then we obtain from ( 3 ) $\Lambda=Q$, i.e., the problem reduces to determining the flux passing through the cube. We first find the magnitude of the flux $Q_{1}$ flowing through the first layer, in which there are $M^{2}$ resistances, while $M=N+1$. We denote by $i=1,2, \ldots, M^{2}$ the numbers ordering the nodes; $\Lambda_{i, i}+M^{2}$, conductivity between neighboring nodes along the $O Z$ axis; $S_{i}$, area per resistance; and $\delta=1 / \mathrm{N}$, thickness of a single layer. The magnitude of the flux $\mathrm{Q}_{1}$ is given by

$$
\begin{equation*}
Q_{1}=\frac{1}{\delta} \sum_{i=1}^{M^{2}} \Lambda_{i, i+M^{2}}\left(t_{i+M^{2}}-t_{i}\right) S_{i}, \tag{4}
\end{equation*}
$$

where $t_{i}+M^{2}$ and $t_{i}$ are the temperatures at neighboring points along the $O Z$ axis.
In order to determine the temperature field of the model, we will write the equation for balance of heat energy at the i-th point:

$$
\begin{gather*}
\left(\Lambda_{i, i+1}+\Lambda_{i, i-1}+\Lambda_{i, i+M}+\Lambda_{i, i-M}+\Lambda_{i, i+M^{2}}+\right. \\
\left.+\Lambda_{i, i-M^{2}}\right) t_{i}-\Lambda_{i-1, i} i_{i-1}-\Lambda_{i+1, i} i_{i+1}-\Lambda_{i-M, i} t_{i-M}-\Lambda_{i+M, i} t_{i+M}-\Lambda_{i-M^{2}, i_{i-M^{2}}-\Lambda_{i+M^{2}, i} t_{i+M^{2}}=0,} \tag{5}
\end{gather*}
$$

where $\Lambda_{i, i-1} ; \Lambda_{i, i+1}$ are the conductivity between neighboring points along the OY axis; $\Lambda_{i, i-M^{2}} ; \Lambda_{i, i+M^{2}}$ are the conductivity between neighboring points along the OZ axis; $\Lambda_{i, i-M} ; \Lambda_{i, i+M}$ are the conductivity between neighboring points along the OX axis. Equation (5) can be written in a different form as follows:

$$
t_{i} \sum_{j=M^{2}+1}^{M^{3}-M^{2}} \Lambda_{i j}-\sum_{i=M^{2}+1}^{M^{3}-M^{2}} t_{j} \Lambda_{i j}=0
$$

and, in addition,

$$
\Lambda_{i j} \neq 0 \text { for }\left\{\begin{array}{l}
j=i-1 ; i+1  \tag{6}\\
j=i-M ; i+M \\
j=i-M^{2} ; i+M^{2}
\end{array}\right.
$$

If Eq. (6) is written for all values of $i=M^{2}+1, M^{2}+2, \ldots, M^{3}-M^{2}$, then we obtain a system of ordinary algebraic equations, whose solution determines the temperature field of the model. According to the conditions of the problem, for a stationary process, the flux passing through a single layer equals the flux passing through the entire cube. The system of equations (4) and (6) was solved using the Monte Carlotechnique on an ES-1022 computer. Using the method indicated, we determined the dimensions of a representative element of a binary heterogeneous system. For this, we calculated the effective conductivity of the $\Lambda$ model, containing from two to five layers. Figure 3a shows the dependence of $\Lambda=f(N, \nu)$ conductivity on the number of layers $N$ for different values of the ratio $\nu=\Lambda_{2} / \Lambda_{1}$. From the results obtained, we can draw the following conclusions:
a) the value of $\Lambda_{2} / \Lambda_{1}$ has practically no effect on the choice of the representative element;
b) the deviation of $\Lambda$ from the limiting value ( $\Lambda_{\infty}$ for $N \rightarrow \infty$ ) increases with increasing concentration of the nonconducting component;
c) over the entire range of concentrations, the results obtained with the five-layer model do not deviate from the limiting values by more than $5 \%$. Thus, the representative element of a binary heterogeneous system must contain at least five layers.

The conductivity of multicomponent systems was subsequently studied using the five-layer model. Figure $3 b$ shows the results of a calculation of $\Lambda$ for a binary system with different values of $\Lambda_{2} / \Lambda_{1}$ over the entire range of variation of concentrations $0 \leq m_{2} \leq 1$ (continuous lines); the dashed lines in this figure indicate the dependences calculated using the equations obtained by simulation, proposed by Dul'nev and Novikov [5, 6].

The conductivity of a three-component system was investigated using the same model. The values of the resistance of the first, second, and third types were distributed among the nodes of the model in two stages. In the first stage, bonds were chosen for which the values of the resistances corresponding to the first component were assigned in a random manner; the number of such bonds was determined by the volume concentration $m_{1}$. At the second stage, the values of the given volume concentration of the remaining two components were renormalized:

$$
\begin{equation*}
m_{2}^{\prime}=\frac{m_{2}}{1-m_{1}}, m_{3}^{\prime}=\frac{m_{3}}{1-m_{1}} \tag{7}
\end{equation*}
$$

where $m_{2}^{\prime}$ and $m_{3}^{\prime}$ are the new values of the volume concentration of the second and third components and, in addition, $m_{2}^{j}+m_{1}^{\prime}=1$. The values of the resistances of the second and third components were distributed according to the values of the concentrations obtained. Figure 4 shows the results of the calculation of the effective conductivity of a three-component system; the dashed lines indicate the values obtained from the model representations based on a self-consistent method [7].

## NOTATION

$\Lambda$, conductivity of the heterogeneous system; $\Lambda_{i}$, conductivity of the i-th component; $m_{i}$, volume concentration of the $i$-th component; $m_{M}$ and $m_{I}$, volume concentrations of the conducting and nonconducting (insulating) components.

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